

2014

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Gudder, S., "Emergence of Four Dimensions in the Causal Set Approach to Discrete Quantum Gravity" (2014). *Mathematics Preprint Series*. 6.

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EMERGENCE OF FOUR DIMENSIONS IN THE CAUSAL SET APPROACH TO DISCRETE QUANTUM GRAVITY

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Abstract

One could begin a study like the present one by simply postulating that our universe is four-dimensional. There are ample reasons for doing this. Experience, observation and experiment all point to the fact that we inhabit a four-dimensional universe. Another approach would be to show that four-dimensions arise naturally from a reasonable model of the universe or multiverse. After reviewing the causal set approach to discrete quantum gravity in Section 1, we shall discuss the emergence of four-dimensions in Section 2. We shall see that certain patterns of four arise that suggest the introduction of a 4-dimensional discrete manifold. In the later sections we shall discuss some consequences of this introduced framework. In particular, we will show that quantum amplitudes can be employed to describe a multiverse dynamics. Moreover, a natural unitary operator together with energy, position and momentum operators will be introduced and their properties studied.

Keywords: Causal sets, discrete quantum gravity, transition amplitudes, four dimensions.

1 Causal Set Approach

We call a finite poset $(x, <)$ a *causet* and interpret $a < b$ in x to mean that b is in the causal future of a [3, 5, 7, 8]. If x and y are causets, we say that x *produces* y (denoted by $x \rightarrow y$) if y is obtained from x by adjoining a single maximal element to x . If $x \rightarrow y$, we call y an *offspring* of x . A *labeling* for a causet x of cardinality $|x|$ is a bijection

$$\ell: x \rightarrow \{1, 2, \dots, |x|\}$$

such that $a, b \in x$ with $a < b$ implies $\ell(a) < \ell(b)$. Two labeled causets x, y are *isomorphic* if there is a *bijection* $\phi: x \rightarrow y$ such that $a < b$ in x if and only if $\phi(a) < \phi(b)$ in y and $\ell[\phi(a)] = \ell(a)$ for all $a \in x$. A causet is *covariant* if it has a unique labeling (up to isomorphisms) and we call a covariant causet a *c-causet* [1, 2, 3, 4]. Denote the set of *c-causets* with cardinality n by \mathcal{P}'_n and the set of all *c-causets* by \mathcal{P}' . It is shown in [2] that any $x \in \mathcal{P}'$ with $x \neq \emptyset$ has a unique producer in \mathcal{P}' and precisely two offspring in \mathcal{P}' . It follows that $|\mathcal{P}'_n| = 2^{n-1}$, $n = 1, 2, \dots$. Two elements $a, b \in x$ are *comparable* if $a < b$ or $b < a$. The *height* $h(a)$ of $a \in x$ is the cardinality, minus one, of the longest path in x that ends with a . It is shown in [2] that a causet x is covariant if and only if $a, b \in x$ are comparable whenever $h(a) \neq h(b)$.

If $x \in \mathcal{P}'$ we call the sets

$$S_j(x) = \{a \in x: h(a) = j\}, \quad j = 0, 1, 2, \dots$$

shells and the sequence of integers $s_i(x) = |S_i(x)|$, $j = 0, 1, \dots$, is the *shell sequence*. A *c-causet* is uniquely determined by its shell sequence and we think of $\{s_i(x)\}$ as describing the “shape” or geometry of x . The tree $(\mathcal{P}', \rightarrow)$ can be thought of as a growth model and an $x \in \mathcal{P}'_n$ is a possible universe at step (time) n . An instantaneous universe $x \in \mathcal{P}'_n$ grows one element at a time in one of two ways. If x has shell sequence $(s_0(x), s_1(x), \dots, s_m(x))$ then $x \rightarrow x_0$ or $x \rightarrow x_1$ where $x_0, x_1 \in \mathcal{P}'_{n+1}$ have shell sequences

$$(s_0(x), s_1(x), \dots, s_m(x) + 1)$$

and

$$(s_0(x), s_1(x), \dots, s_m(x), 1)$$

respectively. In this way, we recursively order the c -causets in \mathcal{P}'_n using the notation $x_{n,j}$, $n = 1, 2, \dots$, $j = 0, 1, \dots, 2^{n-1} - 1$ where $n = |x_{n,j}|$. For example, in terms of their shell sequences we have:

$$\begin{aligned} x_{1,0} &= (1), x_{2,0} = (2), x_{2,1} = (1, 1), x_{3,0} = (3), x_{3,1} = (2, 1), x_{3,2} = (1, 2), x_{3,3} = (1, 1, 1) \\ x_{4,0} &= (4), x_{4,1} = (3, 1), x_{4,2} = (2, 2), x_{4,3} = (2, 1, 1), x_{4,4} = (1, 3), x_{4,5} = (1, 2, 1) \\ x_{4,6} &= (1, 1, 2), x_{4,7} = (1, 1, 1, 1) \end{aligned}$$

In this model, we view a c -causet as a framework or scaffolding for a possible universe. The vertices of x represent small cells that can be empty or occupied by a particle. The shell sequence for x gives the geometry of the framework [3]. Notice that this is a multiverse model in which infinite paths in $(\mathcal{P}', \rightarrow)$ represent the histories of “completed” evolved universes.

We now describe the evolution of a universe as a quantum sequential growth process. In such a process, the probabilities of competing geometries are determined by quantum amplitudes. These amplitudes provide interferences that are characteristic of quantum systems. A *transition amplitude* is a map $\tilde{a}: \mathcal{P}' \times \mathcal{P}' \rightarrow \mathbb{C}$ satisfying $\tilde{a}(x, y) = 0$ if $x \not\rightarrow y$ and $\sum_{y \in \mathcal{P}'} \tilde{a}(x, y) = 1$. Since $x_{n,j}$ only has the offspring $x_{n+1,2j}$ and $x_{n+1,2j+1}$ we have that

$$\sum_{k=0}^1 \tilde{a}(x_{n,j}, x_{n+1,2j+k}) = 1$$

for all $n = 1, 2, \dots$, $j = 0, 1, \dots, 2^{n-1} - 1$. We call \tilde{a} a *unitary transition amplitude* (uta) if \tilde{a} also satisfies $\sum_{y \in \mathcal{P}'} |\tilde{a}(x, y)|^2 = 1$. One might suspect that the restrictions on a uta are so strong that the possibilities are very limited. This would be true if \tilde{a} were real valued. In this case, $\tilde{a}(x, y)$ is 0 or 1. However, in the complex case, the next result, which is proved in [4], shows that there are a continuum of possibilities.

Theorem 1.1. *Two complex numbers a, b satisfy $a + b = |a|^2 + |b|^2 = 1$ if and only if there exists a $\theta \in [0, \pi)$ such that $a = \cos \theta e^{i\theta}$, $b = -i \sin \theta e^{i\theta}$. Moreover, θ is unique.*

If $\tilde{a}: \mathcal{P}' \times \mathcal{P}' \rightarrow \mathbb{C}$ is a uta, we call

$$c_{n,j}^k = \tilde{a}(x_{n,j}, x_{n+1,2j+k}), \quad k = 0, 1$$

the 2-dimensional coupling constants for \tilde{a} . It follows from Theorem 1.1 that there exist $\theta_{n,j} \in [0, \pi)$ such that

$$c_{n,j}^0 = \cos \theta_{n,j} e^{i\theta_{n,j}}, c_{n,j}^1 = -i \sin \theta_{n,j} e^{i\theta_{n,j}}$$

Hence, $c_{n,j}^0 + c_{n,j}^1 = |c_{n,j}^0|^2 + |c_{n,j}^1|^2 = 1$ for all $n = 1, 2, \dots, j = 0, 1, \dots, 2^{n-1} - 1$.

A *path* in \mathcal{P}' is a sequence $\omega = \omega_1 \omega_2 \cdots$ where $\omega_i \in \mathcal{P}'_i$ and $\omega_i \rightarrow \omega_{i+1}$. Similarly, an *n-path* has the form $\omega = \omega_1 \omega_2 \cdots \omega_n$ where again $\omega_i \rightarrow \omega_{i+1}$. We denote the set of paths by Ω' and the set of *n*-paths by Ω'_n . Since every $x \in \mathcal{P}'_n$ has a unique *n*-path terminating at x , we can identify \mathcal{P}'_n with Ω'_n and write $\mathcal{P}'_n \approx \Omega'_n$. If \tilde{a} is a uta and $\omega = \omega_1 \omega_2 \cdots \omega_n \in \Omega'_n$, we define the *amplitude* of ω to be

$$a(\omega) = \tilde{a}(\omega_1, \omega_2) \tilde{a}(\omega_2, \omega_3) \cdots \tilde{a}(\omega_{n-1}, \omega_n)$$

Moreover, we define the *amplitude* of $x \in \mathcal{P}'_n$ to be $a(\omega)$ where ω is the unique path in Ω'_n that terminates at x . For $A \subseteq \mathcal{P}'_n$ we define the *amplitude* of A to be

$$a(A) = \sum \{a(x) : x \in A\}$$

and the *q-measure* of A to be $\mu_n(A) = |a(A)|^2$. We conclude that $\mu_n : 2^{\mathcal{P}'_n} \rightarrow \mathbb{R}^+$ and it is not hard to verify that $\mu(\mathcal{P}'_n) = 1$ [5, 6, 7].

In general, μ_n is not additive so it is not a measure. For this reason $\mu_n(A)$ is interpreted as the quantum propensity of A instead of the quantum probability. Although μ_n is not additive, it satisfies the *grade-2 additivity condition* [5, 6, 7]: if $A, B, C \in 2^{\mathcal{P}'_n}$ are mutually disjoint then

$$\mu_n(A \cup B \cup C) = \mu_n(A \cup B) + \mu_n(A \cup C) + \mu_n(B \cup C) - \mu_n(A) - \mu_n(B) - \mu_n(C)$$

Because of the lack of additivity we have, in general, that

$$\mu_n(\{x, y\}) \neq \mu_n(x) + \mu_n(y) \tag{1.1}$$

for $x, y \in \mathcal{P}'_n$. If (1.1) holds, we say that x and y *interfere* and otherwise we say that x and y *do not interfere*. It is shown in [4] that if x and y have the same producer, then x and y do not interfere.

2 Patterns of Four

This section shows that various patterns of four occur in \mathcal{P}'_n . The *height* $h(x)$ of $x \in \mathcal{P}'$ is the cardinality minus one of the longest paths in x . Equivalently, $h(x)$ is the number of shells minus one in x . For example, since $x_{4,5} = (1, 2, 1)$ we have $h(x_{4,5}) = 2$. The *height sequence* of \mathcal{P}'_n is the sequence of integers

$$(h(x_{n,0}), h(x_{n,1}), \dots, h(x_{n,2^{n-1}-1}))$$

We now display the height sequences of \mathcal{P}'_n for the first few values of n .

$$\mathcal{P}'_1: (0), \mathcal{P}'_2: (0, 1), \mathcal{P}'_3: (0, 1, 1, 2), \mathcal{P}'_4: (0, 1, 1, 2, 1, 2, 2, 3)$$

$$\mathcal{P}'_5: (0, 1, 1, 2, 1, 2, 2, 3, 1, 2, 2, 3, 2, 3, 3, 4)$$

$$\mathcal{P}'_6: (0, 1, 1, 2, 1, 2, 2, 3, 1, 2, 2, 3, 2, 3, 3, 4, 1, 2, 2, 3, 2, 3, 3, 4, 2, 3, 3, 4, 3, 4, 4, 5)$$

Notice the compelling patterns of four that stand out. For example, in \mathcal{P}'_6 we have $(0, 1, 1, 2)$, $(1, 2, 2, 3)$, $(1, 2, 2, 3)$, $(2, 3, 3, 4)$, $(1, 2, 2, 3)$, $(2, 3, 3, 4)$, $(2, 3, 3, 4)$, $(3, 4, 4, 5)$. These patterns have the form $r, r+1, r+1, r+2$. One might also say that there are patterns of two and patterns of eight, but these are not as compelling. We now show why this four pattern occurs. If $x \rightarrow y$ then $h(y) = h(x)$ or $h(x) + 1$. Hence, if \mathcal{P}'_n has height sequence,

$$(r_1, r_2, \dots, r_{2^{n-1}-1})$$

then \mathcal{P}'_{n+1} has height sequence

$$(r_1, r_1 + 1, r_2, r_2 + 1, \dots, r_{2^{n-1}-1}, r_{2^{n-1}-1} + 1)$$

Applying this reasoning again shows that \mathcal{P}'_{n+2} has height sequence

$$(r_1, r_1 + 1, r_1 + 1, r_1 + 2, r_2, r_2 + 1, r_2 + 1, r_2 + 2, \dots, \\ r_{2^{n-1}-1}, r_{2^{n-1}-1} + 1, r_{2^{n-1}-1} + 1, r_{2^{n-1}-1} + 2)$$

As an aside, it is of interest to consider the number of causets $\tau_n(r)$ in \mathcal{P}'_n with height $r = 0, 1, \dots, n-1$. For example, $\tau_5(0) = 1$, $\tau_5(1) = 4$, $\tau_5(2) = 6$, $\tau_5(3) = 4$, $\tau_5(4) = 1$. This suggests the following result.

Theorem 2.1. *For $n = 1, 2, \dots$, $r = 0, 1, \dots, n-1$ we have that $\tau_n(r) = \binom{n-1}{r}$.*

Proof. We employ induction on n . The result $\tau_1(0) = 1$ is clearly true. Assume the result holds for \mathcal{P}'_n and consider causets in \mathcal{P}'_{n+1} with height r . Now any $x \in \mathcal{P}'_n$ with $h(x) = r$ produces a $y \in \mathcal{P}'_n$ with $h(y) = r$. Also any $x \in \mathcal{P}'_n$ with $h(x) = r - 1$ (we can assume $r \neq 0$) produces a $y \in \mathcal{P}'_n$ with $h(y) = r$. By the induction hypothesis

$$\tau_{n+1}(r) = \tau_n(r) + \tau_n(r - 1) = \binom{n-1}{r} + \binom{n-1}{r-1}$$

But a well-known combinatorial identity gives

$$\binom{n-1}{r} + \binom{n-1}{r-1} = \binom{n}{r} \quad \square$$

Corollary 2.2. *The maximum of $\tau_n(r)$ occurs when $r = \lfloor \frac{n-1}{2} \rfloor$ where $\lfloor \cdot \rfloor$ is the floor function.*

We now consider a second example of four patterns. If $x \in \mathcal{P}'$ has shell sequence $(s_0(x), s_1(x), \dots, s_m(x))$, $s_m(x) > 0$ we define $w(x) = s_m(x)$. Thus, $w(x)$ is the cardinality of the highest shell of x . The *width sequence* of \mathcal{P}'_n is the sequence of integers

$$(w(x_{n,0}), w(x_{n,1}), \dots, w(x_{n,2^n-1}))$$

We next display the width sequences of \mathcal{P}'_n for the first few values of n .

$$\mathcal{P}'_1: (1), \mathcal{P}'_2: (2, 1), \mathcal{P}'_3: (3, 1, 2, 1), \mathcal{P}'_4: (4, 1, 2, 1, 3, 1, 2, 1)$$

$$\mathcal{P}'_5: (5, 1, 2, 1, 3, 1, 2, 1, 4, 1, 2, 1, 3, 1, 2, 1)$$

$$\mathcal{P}'_6: (6, 1, 2, 1, 3, 1, 2, 1, 4, 1, 2, 1, 3, 1, 2, 1, 5, 1, 2, 1, 3, 1, 2, 1, 4, 1, 2, 1, 3, 1, 2, 1)$$

As before, the patterns of four $(r, 1, 2, 1)$ stand out clearly.

Our last example of a four pattern stems from quantum mechanics. What we say here applies to any \mathcal{P}'_n for $n \geq 3$ but for simplicity let us consider $\mathcal{P}'_3 = \{x_{3,0}, x_{3,1}, x_{3,2}, x_{3,3}\}$. Since $x_{3,0}$ and $x_{3,1}$ have the same producer, as mentioned in Section 1, $x_{3,0}$ and $x_{3,1}$ do not interfere. Thus, the pair $(x_{3,0}, x_{3,1})$ act classically with no apparent quantum effects. The same applies for the pair $(x_{3,2}, x_{3,3})$. Calling such pairs *siblings*, if we want to include siblings and quantum effects the smallest applicable set is the 4-tuple $(x_{3,0}, x_{3,1}, x_{3,2}, x_{3,3})$. In general, the 4-tuple would have the form

$$(x_{n,4j}, x_{n,4j+1}, x_{n,4j+2}, x_{n,4j+3})$$

There are precisely *four* interfering pairs in such a 4-tuple. These are $(x_{n,4j}, x_{n,4j+2})$, $(x_{n,4j}, x_{n,4j+3})$, $(x_{n,4j+1}, x_{n,4j+2})$ and $(x_{n,4j+1}, x_{n,4j+3})$.

Example. We now show that, in general, $x_{3,0}$ and $x_{3,2}$ interfere. In terms of the coupling constants we have that

$$\mu_3(\{x_{3,0}, x_{3,2}\}) = |a(x_{3,0}) + a(x_{3,2})|^2 = |c_{1,0}^0 c_{2,0}^0 + c_{1,0}^1 c_{2,1}^0|^2 \quad (2.1)$$

On the other hand

$$\mu_3(x_{3,0}) + \mu_3(x_{3,2}) = |a(x_{3,0})|^2 + |a(x_{3,2})|^2 = |c_{1,0}^0|^2 |c_{2,0}^0|^2 + |c_{1,0}^1|^2 |c_{2,1}^0|^2 \quad (2.2)$$

But (2.1) and (2.2) do not agree unless

$$\operatorname{Re}(c_{1,0}^0 c_{2,0}^0 \bar{c}_{1,0}^1 \bar{c}_{2,1}^0) = 0$$

so $x_{3,0}$ and $x_{3,2}$ interfere, in general. The same reasoning also shows that $(x_{3,0}, x_{3,3})$, $(x_{3,1}, x_{3,2})$, $(x_{3,1}, x_{3,3})$ are interfering pairs, in general.

This argument also shows that $x_{3,0}$ and $x_{3,1}$ do not interfere. In this case (2.1) becomes

$$\mu_3(\{x_{3,0}, x_{3,1}\}) = |c_{1,0}^0 c_{2,0}^0 + c_{1,0}^0 c_{2,0}^1|^2 = |c_{1,0}^0|^2 |c_{2,0}^0 + c_{2,0}^1|^2 = |c_{1,0}^0|^2$$

and (2.2) becomes

$$\mu_3(x_{3,0}) + \mu_3(x_{3,1}) = |c_{1,0}^0|^2 (|c_{2,0}^0|^2 + |c_{2,0}^1|^2) = |c_{1,0}^0|^2$$

Hence, (2.1) and (2.2) agree.

3 Four Dimensional Discrete Manifold

The previous section suggests that important patterns occur for c -causet 4-tuples of the form $(x_{n,4j}, x_{n,4j+1}, x_{n,4j+2}, x_{n,4j+3})$. We can write such 4-tuples as $x_{n,4j+k}$, $k = 0, 1, 2, 3$. This indicates that instead of considering the full set of c -causets \mathcal{P}' we should concentrate on the set

$$\mathcal{P} = \cup \{\mathcal{P}'_n : n \text{ odd}\}$$

of odd cardinality c -causets. We then define $\mathcal{P}_n = \mathcal{P}'_{2n-1}$ to be the collection of c -causets with cardinality $2n-1$, $n = 1, 2, \dots$, so that $\mathcal{P} = \cup \mathcal{P}_n$. We now have

$$|\mathcal{P}_n| = |\mathcal{P}'_{2n-1}| = 2^{2n-2} = 4^{n-1}$$

and as before we order the c -causets in \mathcal{P}_n as

$$\mathcal{P}_n = \{x_{n,0}, x_{n,1}, \dots, x_{n,4^{n-1}-1}\}$$

In this case, each $x \in \mathcal{P}$ except $x_{1,0}$ has a unique producer and each $x \in \mathcal{P}$ has four offspring. In particular $x_{n,j} \rightarrow x_{n+1,4j+k}$, $k = 0, 1, 2, 3$, so $(\mathcal{P}, \rightarrow)$ becomes a tree that we interpret as a sequential growth process. The main difference is that $x \rightarrow y$ if y is obtained by first adjoining a maximal element a to x and then adjoining a second maximal element b to $x \cup \{a\}$ so that $y = x \cup \{a, b\}$.

In this framework, \mathcal{P} has the structure of a discrete 4-manifold. A *tangent vector* at $x \in \mathcal{P}$ is a pair (x, y) where $x \in \mathcal{P}_n$, $y \in \mathcal{P}_{n+1}$ and $x \rightarrow y$. Since every $x \in \mathcal{P}$ has four offspring, there are four tangent vectors at x . We denote the tangent vectors at $x_{n,j}$ by $d_{n,j}^k$, $k = 0, 1, 2, 3$ where

$$d_{n,j}^k = (x_{n,j}, x_{n+1,4j+k})$$

We say that two tangent vectors are *incident* if they have the forms (x, y) , (y, z) . As before, an n -*path* in \mathcal{P} is a sequence $\omega = \omega_1 \omega_2 \dots \omega_n$ where $\omega_i \in \mathcal{P}_i$ and $\omega_i \rightarrow \omega_{i+1}$. We denote the set of n -paths by Ω_n . We can consider an n -path as a sequence of tangent vectors

$$\omega = d_{1,0}^{k_1} d_{2,j_2}^{k_2} \dots d_{n-1,j_{n-1}}^{k_{n-1}}$$

where each tangent vector is incident to the next.

A *transition amplitude* $\tilde{a}: \mathcal{P} \times \mathcal{P} \rightarrow \mathbb{C}$ and a *unitary transition amplitude* (uta) are defined as before. Moreover, we call

$$c_{n,j}^k = \tilde{a}(x_{n,j}, x_{n+1,4j+k}), \quad k = 0, 1, 2, 3$$

the *coupling constants* for \tilde{a} . If \tilde{a} is a uta, we have

$$\sum_{k=0}^3 c_{n,j}^k = \sum_{k=0}^3 |c_{n,j}^k| = 1 \quad (3.1)$$

for $n = 1, 2, \dots$, $j = 0, 1, \dots, 4^{n-1} - 1$. As in Section 2, we can identify \mathcal{P}_n with Ω_n and write $\mathcal{P}_n \approx \Omega_n$. If \tilde{a} is a uta and $\omega = \omega_1 \omega_2 \cdots \omega_n \in \Omega_n$ we define the *amplitude* of ω to be

$$a(\omega) = \tilde{a}(\omega_1, \omega_2) \tilde{a}(\omega_2, \omega_3) \cdots \tilde{a}(\omega_{n-1}, \omega_n)$$

Also, we define the *amplitude* of $x \in \mathcal{P}$ to be $a(\omega)$ where ω is the unique path in Ω_n that terminates at x .

Let H_n be the Hilbert space

$$H_n = L_2(\Omega_n) = L_2(\mathcal{P}_n) = \{f: \mathcal{P}_n \rightarrow \mathbb{C}\}$$

with the standard inner product

$$\langle f, g \rangle = \sum_{x \in \mathcal{P}_n} \overline{f(x)} g(x)$$

Let $\hat{x}_{n,j}$ be the unit vector in H_n given by the characteristic function $\chi_{x_{n,j}}$. Then $\dim H_n = 4^{n-1}$ and $\{\hat{x}_{n,j}: j = 0, 1, \dots, 4^{n-1} - 1\}$ forms an orthonormal basis for H_n . For the remainder of this section \tilde{a} is a uta with corresponding coupling constants $c_{n,j}^k$. We now describe the quantum dynamics generated by \tilde{a} . Define the operators $U_n: H_n \rightarrow H_{n+1}$ by

$$U_n \hat{x}_{n,j} = \sum_{k=0}^3 c_{n,j}^k \hat{x}_{n+1,4j+k}$$

and extend U_n to H_n by linearity. The next few theorems generalize results in [4].

Theorem 3.1. (i) *The adjoint of U_n is given by $U_n^*: H_{n+1} \rightarrow H_n$, where*

$$U_n^* \hat{x}_{n+1,4j+k} = \bar{c}_{n,j}^k \hat{x}_{n,j}, \quad k = 0, 1, 2, 3 \quad (3.2)$$

(ii) *U_n is a partial isometry with $U_n^* U_n = I_n$ and*

$$U_n U_n^* = \sum_{j=0}^{4^{n-1}-1} \left| \sum_{k=0}^3 c_{n,j}^k \hat{x}_{n+1,4j+k} \right\rangle \left\langle \sum_{k=0}^3 c_{n,j}^k \hat{x}_{n+1,4j+k} \right| \quad (3.3)$$

Proof. (i) To show that (3.2) holds we have that

$$\begin{aligned} \langle U_n^* \widehat{x}_{n+1,4j'+k'}, \widehat{x}_{n,j} \rangle &= \langle \widehat{x}_{n+1,4j'+k'}, U_n x_{n,j} \rangle \\ &= \left\langle \widehat{x}_{n+1,4j'+k'}, \sum_{k=0}^3 c_{n,j}^k \widehat{x}_{n+1,4j+k} \right\rangle \\ &= c_{n,j'}^{k'} \delta_{j,j'} = \left\langle \overline{c}_{n,j'}^{k'} \widehat{x}_{n,j'}, \widehat{x}_{n,j} \right\rangle \end{aligned}$$

(ii) To show that (3.2) holds we have that

$$U_n^* U_n \widehat{x}_{n,j} = \sum_{k=0}^3 c_{n,j}^k U_n^* \widehat{x}_{n+1,4j+k} = \sum_{k=0}^3 |c_{n,j}^k|^2 \widehat{x}_{n,j} = \widehat{x}_{n,j}$$

Since $\{\widehat{x}_{n,j} : j = 0, 1, \dots, 4^{n-1} - 1\}$ forms an orthonormal basis for H_n , the result follows. Equation (3.3) holds because it is well-known that $U_n U_n^*$ is the projection onto the range of U_n . \square

It follows from Theorem 3.1 that the dynamics $U_n : H_n \rightarrow H_{n+1}$ for a uta \tilde{a} is an isometric operator. As usual a *state* on H_n is a positive operator ρ on H_n with $\text{tr}(\rho) = 1$. A *stochastic state* on H_n is a state ρ that satisfies $\langle \rho 1_n, 1_n \rangle = 1$ where $1_n = \chi_{\mathcal{P}_n}$; that is $1_n(x) = 1$ for all $x \in \mathcal{P}_n$. Notice that $U_n^* 1_{n+1} = 1_n$.

Corollary 3.2. (i) If ρ is a state on H_n , then $U_n \rho U_n^*$ is a state on H_{n+1}
(ii) If ρ is a stochastic state on H_n , then $U_n \rho U_n^*$ is a stochastic state on H_{n+1} .

Corollary 3.2 shows that $\rho \rightarrow U_n \rho U_n^*$ gives a quantum dynamics for states. We now show that a natural stochastic state is generated by a uta \tilde{a} . Since

$$\langle \widehat{x}_{n+1,4j+k}, U_n \widehat{x}_{n,j} \rangle = c_{n,j}^k = \tilde{a}(x_{n,j}, x_{n+1,4j+k})$$

we have for any $\omega = \omega_1 \omega_2 \cdots \omega_n \in \Omega_n$ that

$$a(\omega) = \langle \widehat{\omega}_2, U_1 \widehat{\omega}_1 \rangle \langle \widehat{\omega}_3, U_2 \widehat{\omega}_2 \rangle \cdots \langle \widehat{\omega}_n, U_{n-1} \widehat{\omega}_{n-1} \rangle$$

For $\omega \in \Omega_n$, let $\widehat{\omega} = \chi_{\{\omega\}} \in H_n$. Define the operator ρ_n on H_n by $\langle \widehat{\omega}, \rho_n \widehat{\omega}' \rangle = a(\omega) \overline{a(\omega')}$. Equivalently, we have that $\langle \widehat{x}, \rho_n \widehat{y} \rangle = a(x) \overline{a(y)}$ for every $x, y \in$

\mathcal{P}_n . A straightforward generalization of Theorem 2.4 [4] shows that ρ_n is a stochastic state on H_n . We call ρ_n the *amplitude state* corresponding to \tilde{a} .

We have seen that $x_{n,j} \in \mathcal{P}_n$ produces four offspring $x_{n+1,4j+k} \in \mathcal{P}_{n+1}$, $k = 0, 1, 2, 3$. We call the set

$$(x_{n,j} \rightarrow) = \{x_{n+1,4j+k} : k = 0, 1, 2, 3\} \subseteq \mathcal{P}_{n+1}$$

the *one-step causal future* of $x_{n,j}$. For simplicity we write $(x \rightarrow)$ for the one-step causal future of $x \in \mathcal{P}_n$ and we use the notation $(x \rightarrow)^\wedge = \chi_{(x \rightarrow)}$. A straightforward generalization of Theorem 2.5 [4] shows that the amplitude state sequence ρ_n is *consistent* in the sense that

$$\langle (x \rightarrow)^\wedge, \rho_{n+1}(y \rightarrow)^\wedge \rangle = \langle \hat{x}, \rho_n \hat{y} \rangle$$

for every $x, y \in \mathcal{P}_n$. Consistency is important because it states that the probabilities and propensities given by the dynamics ρ_n are conserved in time.

A vector $v \in H_n$ is a *stochastic state vector* if $\|v\| = \langle v, 1_n \rangle = 1$. We call the vector

$$\hat{a}_n = \sum_{j=0}^{4^n-1} a(x_{n,j}) \hat{x}_{n,j} \in H_n \quad (3.4)$$

an *amplitude vector*. Of course, \hat{a}_n is a stochastic state vector [4]. It is easy to check that $\rho_n = |\hat{a}_n\rangle\langle\hat{a}_n|$ so ρ_n is a pure state. The next result shows that \hat{a}_n has the expected properties.

Theorem 3.3. (i) If $v \in H_n$ is a stochastic state vector then so is $U_n v \in H_{n+1}$. (ii) $U_n \hat{a}_n = \hat{a}_{n+1}$. (iii) $U_n^* \hat{a}_{n+1} = \hat{a}_n$.

Proof. (i) This follows from the fact that U_n is isometric and $U_n^* 1_{n+1} = 1_n$. (ii) This follows from

$$\begin{aligned}
U_n \widehat{a}_n &= U_n \sum_{j=0}^{4^{n-1}-1} a(x_{n,j}) \widehat{x}_{n,j} = \sum_{j=0}^{4^{n-1}-1} a(x_{n,j}) \sum_{k=0}^3 c_{n,j}^k \widehat{x}_{n+1,4j+k} \\
&= \sum_{j=0}^{4^{n-1}-1} \sum_{k=0}^3 a(x_{n,j}) c_{n,j}^k \widehat{x}_{n+1,4j+k} \\
&= \sum_{j=0}^{4^{n-1}-1} \sum_{k=0}^3 a(x_{n+1,4j+k}) \widehat{x}_{n+1,4j+k} = \widehat{a}_{n+1}
\end{aligned}$$

(iii) Applying (ii), we have that

$$U_n^* \widehat{a}_{n+1} = U_n^* U_n \widehat{a}_n = \widehat{a}_n \quad \square$$

4 Coupling Constant Symmetry

In Section 3 we considered 4-dimensional coupling constants $c_{n,j}^k = a(x_{n,4j+k})$, $n = 1, 2, \dots$, $j = 0, 1, \dots, 4^{n-1} - 1$, $k = 0, 1, 2, 3$ that satisfied (3.1). How do we know that such coupling constants exist? Certainly there exist trivial coupling constants satisfying $c_{n,j}^k = 0$ or 1, but what about nontrivial coupling constants? One way to construct such coupling constants is by employing tensor products. This construction does not give all possible coupling constants but it does produce a large number of them and it might have physical significance.

Coupling constants can be thought of as stochastic unit vectors in \mathbb{C}^4 of the form

$$u = (c_{n,j}^0, c_{n,j}^1, c_{n,j}^2, c_{n,j}^3)$$

Such vectors satisfy

$$\langle u, 1_4 \rangle = \|u\| = 1$$

where 1_4 is the vector $(1, 1, 1, 1)$. Now let v_1 and v_2 be stochastic unit vectors in \mathbb{C}^2 . (These are particular types of qubit states.) Theorem 1.1 shows that there are a continuum of such vectors. The tensor product $v_1 \otimes v_2$ is a stochastic unit vector in \mathbb{C}^4 . Indeed, since $1_4 = 1_2 \otimes 1_2$ we have

$$\langle v_1 \otimes v_2, 1_4 \rangle = \langle v_1 \otimes v_2, 1_2 \otimes 1_2 \rangle = \langle v_1, 1_2 \rangle \langle v_2, 1_2 \rangle = 1$$

also

$$\|v_1 \otimes v_2\| = \|v_1\| \|v_2\| = 1$$

Now suppose that $e_{n,j}^k, f_{n,j}^k, k = 0, 1$, are 2-dimensional coupling constants as studied in Section 1. Employing the notation $e_{n,j}^k = (e_{n,j}^0, e_{n,j}^1), f_{n,j}^k = (f_{n,j}^0, f_{n,j}^1)$, let $c_{n,j}^k = e_{n,j}^k \otimes f_{n,j}^k$. We conclude that $c_{n,j}^k$ are 4-dimensional coupling constants that we call *product coupling constants*. To be explicit, we have

$$c_{n,j}^k = \begin{bmatrix} \begin{bmatrix} e_{n,j}^0 \\ e_{n,j}^1 \end{bmatrix} & f_{n,j}^0 \\ \begin{bmatrix} e_{n,j}^0 \\ e_{n,j}^1 \end{bmatrix} & f_{n,j}^1 \end{bmatrix} = \begin{bmatrix} e_{n,j}^0 f_{n,j}^0 \\ e_{n,j}^1 f_{n,j}^0 \\ e_{n,j}^0 f_{n,j}^1 \\ e_{n,j}^1 f_{n,j}^1 \end{bmatrix}$$

Hence, $c_{n,j}^0 = e_{n,j}^0 f_{n,j}^0, c_{n,j}^1 = e_{n,j}^1 f_{n,j}^0, c_{n,j}^2 = e_{n,j}^0 f_{n,j}^1, c_{n,j}^3 = e_{n,j}^1 f_{n,j}^1$. This section shows that product coupling constants have particularly interesting properties.

Let a, b be complex numbers satisfying $a+b = |a|^2 + |b|^2 = 1$. We have seen in Theorem 1.1 that there exists a unique $\theta \in [0, \pi)$ such that $a = \cos \theta e^{i\theta}$. We call θ the *angle for a* and write $\theta = \text{ang}(a)$. We now “double-down” the pair (a, b) to form the unitary matrix.

$$A = \begin{bmatrix} a & b \\ b & a \end{bmatrix}$$

The matrix A is also *stochastic* in the sense that $A1_2 = 1_2$. It follows that $u_1 = 2^{-1/2}1_2$ is a unit eigenvector of A with corresponding eigenvalues 1. It is easy to check that

$$u_2 = 2^{-1/2} = \begin{bmatrix} 1 \\ -1 \end{bmatrix}$$

is the other unit eigenvector of A with corresponding eigenvalue

$$a - b = 2a - 1 = e^{i2\theta}$$

Now let $c + d = |c|^2 + |d|^2 = 1$ and again form the stochastic unitary matrix

$$B = \begin{bmatrix} c & d \\ d & c \end{bmatrix}$$

Letting $\phi = \text{ang}(c)$ we have the following result.

Theorem 4.1. *$A \otimes B$ is a 4×4 stochastic unitary matrix with eigenvalues $1, e^{i2\theta}, e^{i2\phi}, e^{i2(\theta+\phi)}$ and corresponding unit eigenvectors $u_1 \otimes u_1, u_2 \otimes u_1, u_1 \otimes u_2, u_2 \otimes u_2$.*

Proof. $A \otimes B$ is unitary because

$$(A \otimes B)(A \otimes B)^* = (A \otimes B)(A^* \otimes B^*) = (AA^* \otimes BB^*) = I_2 \otimes I_2 = I_4$$

Also, $A \otimes B$ is stochastic because

$$(A \otimes B)1_4 = (A \otimes B)(1_2 \otimes 1_2) = A1_2 \otimes B1_2 = 1_2 \otimes 1_2 = 1_4$$

We have seen that $A \otimes B(u_1 \otimes u_1) = u_1 \otimes u_1$. Moreover,

$$\begin{aligned} A \otimes B(u_2 \otimes u_1) &= Au_2 \otimes Bu_1 = e^{i2\theta}u_2 \otimes u_1 \\ A \otimes B(u_1 \otimes u_2) &= Au_1 \otimes Bu_2 = e^{i2\phi}u_1 \otimes u_2 \\ A \otimes B(u_2 \otimes u_2) &= Au_2 \otimes Bu_2 = e^{i2(\theta+\phi)}u_2 \otimes u_2 \quad \square \end{aligned}$$

Let \tilde{a} be a uta on \mathcal{P} with product coupling constants $c_{n,j}^k, n = 1, 2, \dots, j = 0, 1, \dots, 4^{n-1} - 1, k = 0, 1, 2, 3$. Since $c_{n,j}^k$ are product coupling constants, for fixed n, j we have that $c_{n,j}^k = e_{n,j}^k \otimes f_{n,j}^k$ for 2-dimensional coupling constants $e_{n,j}^k$ and $f_{n,j}^k$. We have seen that

$$A_j = \begin{bmatrix} e_{n,j}^0 & e_{n,j}^1 \\ e_{n,j}^1 & e_{n,j}^0 \end{bmatrix}, \quad B_j = \begin{bmatrix} f_{n,j}^0 & f_{n,j}^1 \\ f_{n,j}^1 & f_{n,j}^0 \end{bmatrix}$$

are stochastic unitary matrices so it follows from Theorem 4.1 that $C_j = A_j \otimes B_j$ is also stochastic and unitary. From the definition of the tensor product we have that

$$C_j = \begin{bmatrix} c_{n,j}^0 & c_{n,j}^1 & c_{n,j}^2 & c_{n,j}^3 \\ c_{n,j}^1 & c_{n,j}^0 & c_{n,j}^3 & c_{n,j}^2 \\ c_{n,j}^2 & c_{n,j}^3 & c_{n,j}^0 & c_{n,j}^1 \\ c_{n,j}^3 & c_{n,j}^2 & c_{n,j}^1 & c_{n,j}^0 \end{bmatrix} \quad (4.1)$$

applying Theorem 4.1 again, we conclude that there exist $\theta_{n,j}^1, \theta_{n,j}^2 \in [0, \pi)$ such that the eigenvalues of C_j are $1, e^{i2\theta_{n,j}^1}, e^{i2\theta_{n,j}^2}, e^{i2(\theta_{n,j}^1 + \theta_{n,j}^2)}$ with corresponding eigenvectors $u_1 \otimes u_1, u_2 \otimes u_1, u_1 \otimes u_2, u_2 \otimes u_2$.

The Hilbert space H_{n+1} can be decomposed into the direct sum

$$H_{n+1} = H_{n+1,0} \oplus H_{n+1,1} \oplus \cdots \oplus H_{n+1,4^{n-1}-1} \quad (4.2)$$

where $\dim H_{n+1,j} = 4$ and an orthonormal basis for $H_{n+1,j}$ is

$$\{\widehat{x}_{n+1,4j+k} : k = 0, 1, 2, 3\} \quad (4.3)$$

We now define a stochastic unitary operator $V_{n+1,j} : H_{n+1,j} \rightarrow H_{n+1,j}$ with matrix representation C_j given by (4.1). To be explicit, we have

$$\begin{bmatrix} V_{n+1,j}(\widehat{x}_{n+1,4j}) \\ V_{n+1,j}(\widehat{x}_{n+1,4j+1}) \\ V_{n+1,j}(\widehat{x}_{n+1,4j+2}) \\ V_{n+1,j}(\widehat{x}_{n+1,4j+3}) \end{bmatrix} = C_j \begin{bmatrix} \widehat{x}_{n+1,4j} \\ \widehat{x}_{n+1,4j+1} \\ \widehat{x}_{n+1,4j+2} \\ \widehat{x}_{n+1,4j+3} \end{bmatrix}$$

We conclude from our previous work that the eigenvalues of $V_{n+1,j}$ are $1, e^{i2\theta_{n,j}^1}, e^{i2\theta_{n,j}^2}, e^{i2(\theta_{n,j}^1 + \theta_{n,j}^2)}$ with corresponding unit eigenvectors

$$\begin{aligned} \widehat{x}_{n+1,j}^0 &= \frac{1}{2}(\widehat{x}_{n+1,4j} + \widehat{x}_{n+1,4j+1} + \widehat{x}_{n+1,4j+2} + \widehat{x}_{n+1,4j+3}) \\ \widehat{x}_{n+1,j}^1 &= \frac{1}{2}(\widehat{x}_{n+1,4j} - \widehat{x}_{n+1,4j+1} + \widehat{x}_{n+1,4j+2} - \widehat{x}_{n+1,4j+3}) \\ \widehat{x}_{n+1,j}^2 &= \frac{1}{2}(\widehat{x}_{n+1,4j} + \widehat{x}_{n+1,4j+1} - \widehat{x}_{n+1,4j+2} - \widehat{x}_{n+1,4j+3}) \\ \widehat{x}_{n+1,j}^3 &= \frac{1}{2}(\widehat{x}_{n+1,4j} - \widehat{x}_{n+1,4j+1} - \widehat{x}_{n+1,4j+2} + \widehat{x}_{n+1,4j+3}) \end{aligned} \quad (4.4)$$

Finally, we define the stochastic unitary operator V_{n+1} on H_{n+1} by

$$V_{n+1} = V_{n+1,0} \oplus V_{n+1,1} \oplus \cdots \oplus V_{n+1,4^{n-1}-1}$$

The eigenvalues of V_{n+1} are 1 (with multiplicity 4^{n-1}), $e^{i2\theta_{n+1,j}^1}, e^{i2\theta_{n+1,j}^2}, e^{i2(\theta_{n+1,j}^1 + \theta_{n+1,j}^2)}$, $j = 0, 1, \dots, 4^{n-1} - 1$. The corresponding eigenvectors are $\widehat{x}_{n+1,j}^k$, $j = 0, 1, \dots, 4^{n-1} - 1$, $k = 0, 1, 2, 3$ given by (4.4).

The operator V_{n+1} provides an intrinsic symmetry on H_{n+1} generated by the coupling constants. We call V_{n+1} the *coupling constant symmetry*. Since V_{n+1} is unitary, it has the form $V_{n+1} = e^{iK_{n+1}}$ where K_{n+1} is a self-adjoint operator called the *coupling energy*. The eigenvectors of K_{n+1} are again $\widehat{x}_{n+1,j}^k$ and the corresponding eigenvalues are 0 (with multiplicity 4^{n-1}), $2\theta_{n+1,j}^1$, $2\theta_{n+1,j}^2$, $2(\theta_{n+1,j}^1 + \theta_{n+1,j}^2)$. Since the eigenvalues of K_{n+1} correspond to energy values, we conclude that the physical significance of the $\theta_{n+1,j}^k$ (and 0) are that they are one-half of energy values.

5 Position and Momentum Observables

It is natural to define the *position observable* Q_n on H_n by $Q_n \widehat{x}_{n,j} = j \widehat{x}_{n,j}$. Then Q_n is a self-adjoint operator with eigenvectors $\widehat{x}_{n,j}$ and corresponding eigenvalues j , $j = 0, 1, \dots, 4^{n-1} - 1$. It is also natural to define the *conjugate momentum observable* P_n on H_n by $P_n = V_n Q_n V_n^*$. Then P_n is a self-adjoint operator with the same eigenvalues and corresponding eigenvectors $V_n \widehat{x}_{n,j}$. Now Q_n describes the total position of a causet $x_{n,j} \in \mathcal{P}_n$ but it is important to describe the coordinate observables Q_n^k , according to the *directions* $k = 0, 1, 2, 3$.

To accomplish this, it is useful to write j in its quartic representation

$$j = j_{n-2} j_{n-3} \cdots j_1 j_0, \quad j_i \in \{0, 1, 2, 3\} \quad (5.1)$$

This representation describes the directions that a path turns when moving from $x_{1,0}$ to $x_{n,j}$. For example (5.1) represents the path that turns in direction j_{n-2} at $x_{1,0}$, then turns in direction j_{n-3}, \dots , and finally turns in direction j_0 just before arriving at $x_{n,j}$. If j has the form (5.1) and $k \in \{0, 1, 2, 3\}$ define

$$j^k = j'_{n-2} j'_{n-3} \cdots j'_1 j'_0$$

where $j'_i = 1$ if $j_i = k$ and $j'_i = 0$ if $j_i \neq k$. Thus, j^k marks the places at which j turns in direction k . Notice that

$$j = \sum_{k=0}^3 k j^k \quad (5.2)$$

For $k = 0, 1, 2, 3$, we define the *coordinate observables* Q_n^k by

$$Q_n^k \widehat{x}_{n,j} = j^k \widehat{x}_{n,j}$$

We see that Q_n^k is a self-adjoint operator whose eigenvalues are the 2^{n-1} possible values $j_{n-1}j_{n-2} \cdots j_1j_0$ where $j_i \in \{0, 1\}$. It follows from (5.2) that

$$Q_n = \sum_{k=0}^3 kQ_n^k \quad (5.3)$$

Example. In H_6 , let $j = 403$ whose quartic representation becomes

$$j = 403 = 256 + 128 + 16 + 3 = 1 \times 4^4 + 2 \times 4^3 + 1 \times 4^2 + 0 \times 4 + 3 = 12103$$

We then have $j^0 = 00010 = 4$, $j^1 = 10100 = 272$, $j^2 = 01000 = 64$, $j^3 = 1$.
We then have

$$0 \cdot j^0 + 1 \cdot j^1 + 2 \cdot j^2 + 3 \cdot j^3 = 272 + 128 + 3 = 403$$

The coordinate observables at $j = 403$ become

$$Q_6^0(\widehat{x}_{6,j}) = 4\widehat{x}_{6,j}, \quad Q_6^1(\widehat{x}_{6,j}) = 272\widehat{x}_{6,j}, \quad Q_6^2(\widehat{x}_{6,j}) = 64\widehat{x}_{6,j}, \quad Q_6^3(\widehat{x}_{6,j}) = \widehat{x}_{6,j}$$

As before, we define the *conjugate momentum observables* P_n^k on H_n by $P_n^k = V_n Q_n^k V_n^*$, $k = 0, 1, 2, 3$. To obtain explicit forms for these operators it is convenient to consider their action on $H_{n+1,j}$ given by (4.2) and having basis (4.3). As in Section 4 we write

$$Q_{n+1}^k = Q_{n+1,0}^k \oplus Q_{n+1,1}^k \oplus \cdots \oplus Q_{n+1,4^{n-1}-1}^k$$

and

$$P_{n+1}^k = P_{n+1,0}^k \oplus P_{n+1,1}^k \oplus \cdots \oplus P_{n+1,4^{n-1}-1}^k$$

Now $x_{n,j} \rightarrow x_{n+1,4j}, x_{n+1,4j+1}, x_{n+1,4j+2}, x_{n+1,4j+3}$ and in quartic notation we have $j = j_{n-2}j_{n-3} \cdots j_1j_0$ while for $k = 0, 1, 2, 3$ we have

$$4j + k = j_{n-2}j_{n-3} \cdots j_1j_0k$$

It follows that

$$(4j + \ell)^k = \begin{cases} j^k 1 & \text{if } \ell = k \\ j^k 0 & \text{if } \ell \neq k \end{cases}$$

Hence,

$$Q_{n+1}^k \widehat{x}_{n+1,4j+\ell} = (4j + \ell)^k \widehat{x}_{n+1,4j+\ell} = \begin{cases} j^k 1 \widehat{x}_{n+1,4j+\ell} & \text{if } \ell = k \\ j^k 0 \widehat{x}_{n+1,4j+\ell} & \text{if } \ell \neq k \end{cases}$$

Thus, in the standard basis $\{\widehat{x}_{n+1,4j+k} : k = 0, 1, 2, 3\}$ for $H_{n+1,j}$, the operator $Q_{n+1,j}^0$ has matrix form

$$\begin{aligned} Q_{n+1,j}^0 &= \begin{bmatrix} j^0 1 & 0 & 0 & 0 \\ 0 & j^0 0 & 0 & 0 \\ 0 & 0 & j^0 0 & 0 \\ 0 & 0 & 0 & j^0 0 \end{bmatrix} = (j^0 0)I + \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \\ &= (j^0 0)I + |\widehat{x}_{n+1,4j}\rangle\langle\widehat{x}_{n+1,4j}| \end{aligned}$$

Similarly, we have for $k = 1, 2, 3$ that

$$Q_{n+1,j}^k = (j^k 0)I + |\widehat{x}_{n+1,4j+k}\rangle\langle\widehat{x}_{n+1,4j+k}| \quad (5.4)$$

It follows that the operators $P_{n+1,j}^k$ have the form

$$\begin{aligned} P_{n+1,j}^k &= V_{n+1,j} Q_{n+1,j}^k V_{n+1,j}^* \\ &= (j^k 0)I + |V_{n+1}\widehat{x}_{n+1,4j+k}\rangle\langle V_{n+1}\widehat{x}_{n+1,4j+k}| \end{aligned} \quad (5.5)$$

for $k = 0, 1, 2, 3$

As usual in quantum theory, it is of interest to find the commutator of $Q_{n+1,j}^k$ and $P_{n+1,j}^\ell$. This becomes

$$\begin{aligned} [Q_{n+1,j}^k, P_{n+1,j}^\ell] &= Q_{n+1,j}^k P_{n+1,j}^\ell - P_{n+1,j}^\ell Q_{n+1,j}^k \\ &= |\widehat{x}_{n+1,4j+k}\rangle\langle\widehat{x}_{n+1,4j+k}||V_{n+1}\widehat{x}_{n+1,4j+\ell}\rangle\langle V_{n+1}\widehat{x}_{n+1,4j+\ell}| \\ &\quad - |V_{n+1}\widehat{x}_{n+1,4j+\ell}\rangle\langle V_{n+1}\widehat{x}_{n+1,4j+\ell}||\widehat{x}_{n+1,4j+k}\rangle\langle\widehat{x}_{n+1,4j+k}| \end{aligned}$$

Letting

$$R_{n+1,j}^{k,\ell} = \langle \widehat{x}_{n+1,4j+k}, V_{n+1} \widehat{x}_{n+1,4j+\ell} | \widehat{x}_{n+1,4j+k} \rangle \langle V_{n+1} \widehat{x}_{n+1,4j+\ell} |$$

we have that

$$[Q_{n+1,j}^k, P_{n+1,j}^\ell] = R_{n+1,j}^{k,\ell} - (R_{n+1,j}^{k,\ell})^* = \frac{1}{2i} \text{Im}(R_{n+1,j}^{k,\ell}) \quad (5.6)$$

We can apply (4.1) and (5.6) to find explicit expressions for the commutators. For example

$$\begin{aligned} R_{m+1,k}^{0,0} &= \langle \widehat{x}_{n+1,4j}, V_{n+1} \widehat{x}_{n+1,4j} | \widehat{x}_{n+1,4j} \rangle \langle V_{n+1} \widehat{x}_{n+1,4j} | \\ &= c_{n,j}^0 \left[|\widehat{x}_{n+1,4j} \rangle \left\langle \sum_{k=0}^3 c_{n,j}^k \widehat{x}_{n+1,4j+k} \right. \right] \\ &= c_{n,j}^0 \left[\sum_{k=0}^3 \bar{c}_{n,j}^k |\widehat{x}_{n+1,4j} \rangle \langle \widehat{x}_{n+1,4j+k} | \right] \end{aligned}$$

The matrix representation for $R_{n+1,j}^{0,0}$ is

$$R_{n+1,j}^{0,0} = c_{n,j}^0 \begin{bmatrix} \bar{c}_{n,j}^0 & \bar{c}_{n,j}^1 & \bar{c}_{n,j}^2 & \bar{c}_{n,j}^3 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

We conclude that

$$\begin{aligned} [Q_{n+1,j}^0, P_{n+1,j}^0] &= R_{n+1,j}^{0,0} - (R_{n+1,j}^{0,0})^* \\ &= \begin{bmatrix} 0 & c_{n,j}^0 \bar{c}_{n,j}^1 & c_{n,j}^0 \bar{c}_{n,j}^2 & c_{n,j}^0 \bar{c}_{n,j}^3 \\ -\bar{c}_{n,j}^0 c_{n,j}^1 & 0 & 0 & 0 \\ -\bar{c}_{n,j}^0 c_{n,j}^2 & 0 & 0 & 0 \\ -\bar{c}_{n,j}^0 c_{n,j}^3 & 0 & 0 & 0 \end{bmatrix} \quad (5.7) \end{aligned}$$

As another example we have

$$\begin{aligned}
[Q_{n+1,j}^1, P_{n+1,j}^0] &= R_{n+1,j}^{1,0} - (R_{n+1,j}^{1,0})^* \\
&= \begin{bmatrix} 0 & -\bar{c}_{n,j}^1 c_{n,j}^0 & 0 & 0 \\ c_{n,j}^1 \bar{c}_{n,j}^0 & 0 & c_{n,j}^1 \bar{c}_{n,j}^2 & c_{n,j}^1 \bar{c}_{n,j}^3 \\ 0 & -\bar{c}_{n,j}^1 c_{n,j}^2 & 0 & 0 \\ 0 & -\bar{c}_{n,j}^1 c_{n,j}^3 & 0 & 0 \end{bmatrix}
\end{aligned}$$

The other commutators are similar and the full operators are given by

$$[Q_{n+1}^k, P_{n+1}^\ell] = [Q_{n+1,0}^k, P_{n+1,0}^\ell] \oplus \cdots \oplus [Q_{n+1,4^{n-1}-1}^k, P_{n+1,4^{n-1}-1}^\ell]$$

Any one of these commutators A is anti-self-adjoint in the sense that $A^* = -A$. Thus they have a complete set of eigenvectors and corresponding eigenvalues. The next result considers the case of $[Q_{n+1,j}^0, P_{n+1,j}^0]$ and the others are similar. We assume that $c_{n,j}^0 \neq 0$ because otherwise the result is trivial.

Theorem 5.1. *The eigenvalues of $[Q_{n+1,j}^0, P_{n+1,j}^0]$ are 0 (with multiplicity 2) and*

$$\lambda_{\pm} = \pm i |c_{n,j}^0| \left[1 - |c_{n,j}^0|^2\right]^{1/2} \text{ with corresponding eigenvectors}$$

$$v_r = (0, a_r, b_r, c_r), \quad r = 1, 2$$

where $(a_r, b_r, c_r) \perp (\bar{c}_{n,j}^1, \bar{c}_{n,j}^2, \bar{c}_{n,j}^3)$, $r = 1, 2$ with $v_1 \perp v_2$, $v_r \neq 0$, and

$$v_{\pm} = (\lambda_{\pm}, -\bar{c}_{n,j}^0 c_{n,j}^1, -\bar{c}_{n,j}^0 c_{n,j}^2, -\bar{c}_{n,j}^0 c_{n,j}^3)$$

Proof. From its form in (5.7) it is clear that $[Q_{n+1,j}^0, P_{n+1,j}^0] v_r = 0$, $r = 1, 2$. Moreover, from (5.7) we have

$$\begin{aligned}
[Q_{n+1,j}^1, P_{n+1,j}^0] v_{\pm} &= \begin{bmatrix} -|c_{n,j}^0|^2 (1 - |c_{n,j}^0|^2) \\ \lambda_{\pm}(-\bar{c}_{n,j}^0 c_{n,j}^1) \\ \lambda_{\pm}(-\bar{c}_{n,j}^0 c_{n,j}^2) \\ \lambda_{\pm}(-\bar{c}_{n,j}^0 c_{n,j}^3) \end{bmatrix} \\
&= \lambda_{\pm} \begin{bmatrix} -|c_{n,j}^0|^2 (1 - |c_{n,j}^0|^2) / \lambda_{\pm} \\ -\bar{c}_{n,j}^0 c_{n,j}^1 \\ -\bar{c}_{n,j}^0 c_{n,j}^2 \\ -\bar{c}_{n,j}^0 c_{n,j}^3 \end{bmatrix} \\
&= \lambda_{\pm} v_{\pm} \quad \square
\end{aligned}$$

Of course, we can now apply Theorem 5.1 to find the eigenvalues and eigenvectors of $[Q_{n+1}^0, P_{n+1}^0]$.

6 Expectations

In this section we compute expectations of observables considered in Section 5. These expectations will be computed relative to the amplitude state \widehat{a}_n of Equation (3.4). We denote the expectation of an operator A in the state \widehat{a}_n by

$$E_n(A) = \langle \widehat{a}_n, A \widehat{a}_n \rangle$$

For the position observable Q_n , we have that

$$\begin{aligned}
E_n(Q_n) &= \langle \widehat{a}_n, Q_n \widehat{a}_n \rangle = \sum_{r=0}^{4^{n-1}-1} \sum_{s=0}^{4^{n-1}-1} \bar{a}(x_{n,r}) s a(x_{n,s}) \delta_{r,s} \\
&= \sum_{j=0}^{4^{n-1}-1} j |a(x_{n,j})|^2
\end{aligned}$$

which is not at all surprising. In a similar way, we obtain

$$E_n(Q_n^k) = \sum_{j=0}^{4^{n-1}-1} j^k |a(x_{n,j})|^2$$

To describe momentum operators, it is again convenient to consider H_{n+1} with the amplitude state

$$\widehat{a}_{n+1} = \sum_{k=0}^3 \sum_{j=0}^{4^{n-1}-1} a(x_{n+1,4j+k}) \widehat{x}_{n+1,4j+k}$$

As before, we first consider $P_{n+1,j}^k$ on $H_{n+1,j}$ and define

$$\widehat{a}_{n+1,j} = \sum_{k=0}^3 a(x_{n+1,4j+k}) \widehat{x}_{n+1,4j+k}$$

By (5.5) we have that

$$\begin{aligned} E_{n+1,j}(P_{n+1,j}^k) &= \langle \widehat{a}_{n+1,j}, P_{n+1,j} \widehat{a}_{n+1,j} \rangle \\ &= (j^k 0) \|\widehat{a}_{n+1,j}\|^2 + |\langle V_{n+1} \widehat{x}_{n+1,4j+k}, \widehat{a}_{n+1,j} \rangle|^2 \end{aligned}$$

In particular,

$$\begin{aligned} E_{n+1,j}(P_{n+1,j}^0) &= (j^0 0) \|\widehat{a}_{n+1,j}\|^2 + |\langle V_{n+1} \widehat{x}_{n+1,4j}, \widehat{a}_{n+1,j} \rangle|^2 \\ &= (j^0 0) \|\widehat{a}_{n+1,j}\|^2 + \left| \left\langle \sum_{k=0}^3 c_{n,j}^k \widehat{x}_{n+1,4j+k}, \widehat{a}_{n+1,j} \right\rangle \right|^2 \\ &= (j^0 0) \|\widehat{a}_{n+1,j}\|^2 + \left| \sum_{k=0}^3 \bar{c}_{n,j}^k a(x_{n+1,4j+k}) \right|^2 \\ &= (j^0 0) \|\widehat{a}_{n+1,j}\|^2 + \left| \sum_{k=0}^3 \bar{c}_{n,j}^k c_{n,j}^k a(x_{n,j}) \right|^2 \\ &= (j^0 0) \|\widehat{a}_{n+1,j}\|^2 + |a(x_{n,j})|^2 \\ &= [(j^0 0) + 1] |a(x_{n,j})|^2 \end{aligned}$$

We conclude that

$$E_{n+1}(P_{n+1}^0) = \sum_{j=0}^{4^{n-1}-1} (j^0 0) |a(x_{n,j})|^2 + 1$$

Surprisingly, we obtain a slightly different expression for $E_{n+1,j}(P_{n+1}^1)$. In particular,

$$E_{n+1,j}(P_{n+1}^1) = (j^1 0) \|\widehat{a}_{n+1,j}\|^2 + |\langle V_{n+1} \widehat{x}_{n+1,4j+1}, \widehat{a}_{n+1,j} \rangle|^2$$

But

$$V_{n+1} \widehat{x}_{n+1,4j+1} = c_{n,j}^1 \widehat{x}_{n+1,4j} + c_{n,j}^0 \widehat{x}_{n+1,4j+1} + c_{n,j}^3 \widehat{x}_{n+1,4j+2} + c_{n,j}^2 \widehat{x}_{n+1,4j+3}$$

It follows that

$$\begin{aligned} \langle V_{n+1} \widehat{x}_{n+1,4j+1}, \widehat{a}_{n+1,j} \rangle &= \bar{c}_{n,j}^1 a(x_{n+1,4j}) + \bar{c}_{n,j}^0 a(x_{n+1,4j+1}) + \bar{c}_{n,j}^3 a(x_{n+1,4j+2}) \\ &\quad + \bar{c}_{n,j}^2 a(x_{n+1,4j+3}) \\ &= (\bar{c}_{n,j}^1 c_{n,j}^0 + \bar{c}_{n,j}^0 c_{n,j}^1 + \bar{c}_{n,j}^3 c_{n,j}^2 + \bar{c}_{n,j}^2 c_{n,j}^3) a(x_{n,j}) \\ &= 0 \end{aligned}$$

Hence,

$$E_{n+1,j}(P_{n+1}^1) = (j^1 0) |a(x_{n,j})|^2$$

so that

$$E_{n+1}(P_{n+1}^1) = \sum_{j=0}^{4^{n-1}-1} (j^1 0) |a(x_{n,j})|^2$$

In a similar way we have

$$E_{n+1}(P_{n+1}^k) = \sum_{j=0}^{4^{n-1}-1} (j^k 0) |a(x_{n,j})|^2$$

for $k = 2, 3$.

Even though $[Q_{n+1}^k, P_{n+1}^\ell]$ is not self-adjoint, the operator $i [Q_{n+1}^k, P_{n+1}^\ell]$ is self-adjoint and it is of interest to find its expectation. We begin with

$$E_{n+1,j} \{ [Q_{n+1,j}^0, P_{n+1,j}^0] \} = \langle \widehat{a}_{n+1,j}, [Q_{n+1,j}^0, P_{n+1,j}^0] \widehat{a}_{n+1,j} \rangle$$

Applying (5.7) we have that

$$\begin{aligned}
[Q_{n+1,j}^0, P_{n+1,j}^0] \hat{a}_{n+1,j} &= \begin{bmatrix} c_{n,j}^0 \sum_{k=1}^3 \bar{c}_{n,j}^k a(x_{n+1,4j+k}) \\ -\bar{c}_{n,j}^0 c_{n,j}^1 a(x_{n+1,4j}) \\ -\bar{c}_{n,j}^0 c_{n,j}^2 a(x_{n+1,4j}) \\ -\bar{c}_{n,j}^0 c_{n,j}^3 a(x_{n+1,4j}) \end{bmatrix} \\
&= a(x_{n,j}) \begin{bmatrix} c_{n,j}^0 \sum_{k=1}^3 |c_{n,j}^k|^2 \\ -|c_{n,j}^0|^2 c_{n,j}^1 \\ -|c_{n,j}^0|^2 c_{n,j}^2 \\ -|c_{n,j}^0|^2 c_{n,j}^3 \end{bmatrix}
\end{aligned}$$

We conclude that

$$\begin{aligned}
&E_{n+1,j} \{ [Q_{n+1,j}^0, P_{n+1,j}^0] \} \\
&= |a(x_{n,j})|^2 \left[|c_{n,j}^0|^2 \sum_{k=1}^3 |c_{n,j}^k|^2 - |c_{n,j}^0|^2 |c_{n,j}^1|^2 - |c_{n,j}^0|^2 |c_{n,j}^2|^2 - |c_{n,j}^0|^2 |c_{n,j}^3|^2 \right] \\
&= 0
\end{aligned}$$

It follows that $E_{n+1} \{ [Q_{n+1}^0, P_{n+1}^0] \} = 0$. Hence, for the amplitude state there is no lower bound for the product of the variances of Q_{n+1}^0 and P_{n+1}^0 as in the Heisenberg uncertainty relation. In a similar way, one can show that

$$E_{n+1} \{ [Q_{n+1}^k, P_{n+1}^\ell] \} = 0$$

for $k, \ell = 0, 1, 2, 3$.

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